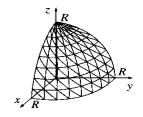
九十四學年第一學期 PHYS2310 電磁學 期中考試題(共兩頁)

[Griffiths Ch. 1-3] 2005/11/15, 10:10am-12:00am, 教師:張存續

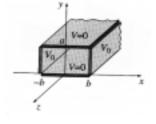
記得寫上學號,班別及姓名等。請依題號順序每頁答一題。

- 1. (4%,4%,4%) Show that
 - (a) $\delta(-x) = \delta(x)$,
 - (b) $\delta(f(x)) = \sum_{i} \frac{1}{\left| \frac{df}{dx}(x_i) \right|} \delta(x x_i)$ where f(x) is assumed to have only simple zero, located at $x = x_i$
 - (c) $\delta(x^2 a^2) = \frac{1}{2|a|} [\delta(x a) + \delta(x + a)]$ [Hint: use the result of (b)]
- 2. (4%, 6%) A vector function $\mathbf{v} = r^2 \cos \theta \hat{\mathbf{r}} + r^2 \cos \phi \hat{\mathbf{\theta}} r^2 \cos \theta \sin \phi \hat{\mathbf{\phi}}$.
 - (a) Find $\nabla \cdot \mathbf{v}$ in spherical coordinate
 - (b) Check the divergence theorem, $\oint \mathbf{v} \cdot d\mathbf{a} = \int_{v} (\nabla \cdot \mathbf{v}) d\tau$, using as your volume one octant of the sphere of radius R. Write down all the surface integrals explicitly.



- 3. (4%, 4%, 4%, 4%) A metal sphere of radius R, carrying charge q, is surrounded by a thick concentric metal shell (inner radius a, outer radius b). The shell carries no net charge.
 - (a) Find the surface charge density σ at R, at a, and at b.
 - (b) Find the electric field at following two regions $R \le r \le a$ and $r \ge b$.
 - (c) Find the potential at the center (r = 0), using the infinity as the reference point.
 - (d) If the outer shell is grounded, what is the potential of the inner sphere and what is the new surface charge density at b, $\sigma(b)$?
- 4. (4%, 4%, 4%)
 - (a) Prove the normal component of **E** is discontinuous at any boundary, using Divergence theorem.
 - (b) Prove the tangential component of E is always continuous, using Stoke's theorem.
 - (c) Write down the normal and tangential component of electric fields immediately outside a metal surface with surface charge density σ .

- 5. (4%, 4%, 4%) The ratio of the magnitude of the charge on one conductor to the magnitude of the potential difference is called the capacitance, C. Find the capacitance of
 - (a) two large, flat, conducting sheets of area A, separated by a small distance d;
 - (b) two concentric conducting spheres with radii a, b (b > a);
 - (c) two concentric conducting cylinders of length L, large compared to their radii a, b (b>a);
- 6. (7%, 7%) Separation of variable in spherical coordinate:
 - (a) Two concentric conducting shells with radii a, b (b>a). The inner shell is connected to a potential of V_0 , while the outer shell is grounded. Find the potential at a < r < b, and r>b.
 - (b) The potential $V(R_0, \theta) = V_0 \sin^2 \theta$ is specified on the surface of a metal sphere, of radius R_0 . Find the potential outside the sphere. [Hint: use Legentre polynomials] $P_0(x) = 0$, $P_1(x) = x$, and $P_2(x) = (3x^2 1)/2$
- 7. (4%, 5%, 5%) Two infinitely long grounded metal plates, again at y=0 and y=a, are connected at $x=\pm b$ by metal strips maintained at a constant potential V_0 , as shown in the figure.
 - (a) Write down the boundary conditions.
 - (b) Write down the general solutions.
 - (c) Find the potential inside the pipe.



- 8. (4%, 3%, 3%) An idea electric dipole **p** is situated at the origin, and points in the z direction. An electric charge q, of mass m, is released from rest at a point in the xy plane. The potential of the dipole is $V(\mathbf{r}) = (1/4\pi\varepsilon_0)(p\cos\theta/r^2)$ and the gravitational force points in the -z direction.
 - (a) Find the electric force between the dipole and the charge.
 - (b) Find the total force (electrical and gravitational) on the charge.
 - (c) Find the electrical potential energy.



1.

(a)
$$\int_{-\infty}^{\infty} f(x)\delta(-x)dx = -\int_{-\infty}^{-\infty} f(-y)\delta(y)dy = \int_{-\infty}^{\infty} f(-y)\delta(y)dy = f(0)$$
$$\int_{-\infty}^{\infty} f(x)\delta(x)dx = f(0) \qquad \therefore \quad \delta(-x) = \delta(x)$$

(b)

$$\int_{-\infty}^{\infty} g(x)\delta(f(x))dx = \sum_{i} \int_{x_{i}-\varepsilon}^{x_{i}+\varepsilon} g(x)\delta(\frac{df}{dx}\Big|_{x_{i}} (x-x_{i}))dx$$

$$\therefore \int_{-\infty}^{\infty} g(x)\delta(k(x-x_{i}))dx = \frac{1}{|k|}g(x_{i})$$

$$\Rightarrow \int_{x_{i}-\varepsilon}^{x_{i}+\varepsilon} g(x)\delta(\frac{df}{dx}\Big|_{x_{i}} (x-x_{i}))dx = \frac{1}{\left|\frac{df}{dx}\Big|_{x_{i}}\right|}g(x_{i})$$

$$\therefore \int_{-\infty}^{\infty} g(x)\delta(f(x))dx = \sum_{i} \frac{1}{\left|\frac{df}{dx}\Big|_{x_{i}}\right|}g(x_{i}) \Rightarrow \delta(f(x)) = \sum_{i} \frac{1}{\left|\frac{df}{dx}\Big|_{x_{i}}\right|}\delta(x-x_{i})$$

(c) Let
$$f(x) = x^2 - a^2$$
, $\delta(f(x)) = \sum_{i} \frac{1}{\left| \frac{df}{dx} \right|_{x_i}} \delta(x - x_i) = \frac{1}{|2a|} [\delta(x - a) + \delta(x + a)]$

2.

(a)
$$\mathbf{v} = r^{2} \cos \theta \hat{\mathbf{r}} + r^{2} \cos \phi \hat{\mathbf{\theta}} - r^{2} \cos \theta \sin \phi \hat{\mathbf{\phi}}$$

$$\nabla \cdot \mathbf{v} = \frac{1}{r^{2}} \frac{\partial}{\partial r} (r^{2} r^{2} \cos \theta) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta r^{2} \cos \phi) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} (-r^{2} \cos \theta \sin \phi)$$

$$= 4r \cos \theta + r \frac{\cos \theta}{\sin \theta} \cos \phi - r \frac{\cos \theta}{\sin \theta} \cos \phi = 4r \cos \theta$$

(b)
$$\int_{v} (\nabla \cdot \mathbf{v}) d\tau = \int_{v} 4r \cos \theta d\tau = \int_{0}^{R} \int_{0}^{\pi/2} \int_{0}^{\pi/2} 4r \cos \theta r^{2} \sin \theta dr d\theta d\phi = R^{4} \frac{1}{2} \frac{\pi}{2} = \frac{\pi R^{4}}{4}$$

$$\oint \mathbf{v} \cdot d\mathbf{a} = \int_{S_{1}} R^{2} \cos \theta R^{2} \sin \theta d\theta d\phi \left(= \frac{\pi R^{4}}{4} \right) + \int_{S_{xy}@\theta = \frac{\pi}{2}} (r^{2} \cos \phi) r dr d\phi \left(= \frac{1}{4} R^{4} \right)$$

$$+ \int_{S_{xx}@\phi = 0} (-r^{2} \cos \theta \sin 0) r dr d\theta \left(= 0 \right) + \int_{S_{yz}@\phi = \frac{\pi}{2}} (-r^{2} \cos \theta \sin \frac{\pi}{2}) r dr d\theta \left(= -\frac{1}{4} R^{4} \right)$$

$$= \frac{\pi R^{4}}{4}$$

3.

(a)
$$\sigma(R) = \frac{q}{4\pi R^2}$$
, $\sigma(a) = \frac{-q}{4\pi a^2}$, and $\sigma(b) = \frac{q}{4\pi b^2}$

(b) Use Gauss's law, we obtain $E = \frac{q}{4\pi\varepsilon_0 r^2}$ for both two regions, $R \le r \le a$ and $r \ge b$.

(c)
$$V(r) = -\int_{\infty}^{r} \mathbf{E} \cdot d\mathbf{r}$$
. $V(0) = V(R) = -\int_{\infty}^{b} \mathbf{E} \cdot d\mathbf{r} - \int_{a}^{R} \mathbf{E} \cdot d\mathbf{r} = \frac{q}{4\pi\varepsilon_{0}} \left(\frac{1}{b} + \frac{1}{R} - \frac{1}{a}\right)$

(d)
$$V(r) = -\int_{a}^{r} \mathbf{E} \cdot d\mathbf{r}$$
, $V(R) = -\int_{a}^{R} \mathbf{E} \cdot d\mathbf{r} = \frac{q}{4\pi\varepsilon_{0}} (\frac{1}{R} - \frac{1}{a})$
 $\sigma(b) = 0$

4.

(a) Consider a Gaussian pillbox. Gauss's law states that $\oint_{S} \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{enc}}{\varepsilon_{0}} = \frac{\sigma A}{\varepsilon_{0}}$

The sides of the pillbox contribute nothing to the flux, in the limit as the thickness ε goes to

zero.
$$(E_{\text{above}}^{\perp} - E_{\text{below}}^{\perp})A = \frac{\sigma A}{\varepsilon_0} \implies (E_{\text{above}}^{\perp} - E_{\text{below}}^{\perp}) = \frac{\sigma}{\varepsilon_0}$$

(b) Consider a thin rectangular loop. The curl of the electric field states that $\oint_P \mathbf{E} \cdot d\ell = 0$

The ends gives nothing (as $\rightarrow \epsilon 0$), and the sides give $(E_{\text{above}}^{"} - E_{\text{below}}^{"})\ell = 0 \implies E_{\text{above}}^{"} = E_{\text{below}}^{"}$

(c) Inside a metal, the electric field is zero
$$\mathbf{E}_{\text{below}} = 0$$
, so $\mathbf{E}_{\text{above}} = \frac{\sigma}{\varepsilon_0} \hat{\mathbf{n}}$ \Rightarrow
$$\begin{cases} E_{above}^{\perp} = \frac{\sigma}{\varepsilon_0} \\ E_{above}^{\parallel} = 0 \end{cases}$$

5.

(a)
$$V = Ed = \frac{Q}{A\varepsilon_0}d$$
, $C = \frac{Q}{V} = \frac{\varepsilon_0 A}{d}$

(b)
$$V = -\int_a^b \mathbf{E} \cdot d\mathbf{r} = \frac{Q}{4\pi\varepsilon_0} (\frac{1}{a} - \frac{1}{b}), \quad C = \frac{Q}{V} = 4\pi\varepsilon_0 \frac{ab}{b-a}$$

(c)
$$V = -\int_a^b \mathbf{E} \cdot d\mathbf{r} = \frac{Q}{2\pi\varepsilon_0 L} \ln \frac{b}{a}$$
, $C = \frac{Q}{V} = 2\pi\varepsilon_0 L / \ln \frac{b}{a}$

6.

(a)



Boundary condition
$$\begin{cases} (i) \ V(a) = V_0 \\ (ii) \ V(b) = 0 \\ (iii) \ V(\infty) = 0 \end{cases}$$

General solution
$$V(r,\theta) = \sum_{\ell=0}^{\infty} (A_{\ell} r^{\ell} + B_{\ell} r^{-(\ell+1)}) P_{\ell}(\cos \theta) = A_0 + B_0 \frac{1}{r}$$

a < r < b

$$\begin{cases} \text{B.C. (i)} \to A_0 + B_0 \frac{1}{a} = V_0 \\ \text{B.C. (ii)} \to A_0 + B_0 \frac{1}{b} = 0 \end{cases} \Rightarrow \begin{cases} B_0 = V_0 \frac{ab}{b-a} \\ A_0 = -V_0 \frac{a}{b-a} \end{cases} \therefore V(r) = \frac{aV_0}{b-a} (-1 + \frac{b}{r})$$

r > b

$$\begin{cases} \text{B.C. (ii)} \rightarrow A_0 + B_0 \frac{1}{b} = 0 \\ \text{B.C. (iii)} \rightarrow A_0 = 0 \end{cases} \Rightarrow \begin{cases} A_0 = 0 \\ B_0 = 0 \end{cases} \therefore V(r) = 0$$

(b)

Boundary condition
$$\begin{cases} (i) \ V(R_0, \theta) = V_0 \sin^2 \theta \\ (ii) \lim_{r \to \infty} V(r, \theta) = 0 \end{cases}$$

General solution $V(r,\theta) = \sum_{\ell=0}^{\infty} (A_{\ell}r^{\ell} + B_{\ell}r^{-(\ell+1)})P_{\ell}(\cos\theta)$

B.C. (i)
$$\rightarrow A_{\ell} = 0$$

B.C. (ii)
$$\rightarrow B_0 R_0^{-1} + B_1 R_0^{-2} \cos \theta + B_2 R_0^{-3} (\frac{3}{2} \cos^2 \theta - \frac{1}{2}) = V_0 (1 - \cos^2 \theta)$$

$$\begin{cases} B_0 R_0^{-1} - \frac{1}{2} B_2 R_0^{-3} = V_0 \\ \frac{3}{2} B_2 R_0^{-3} = -V_0 \\ B_1 = 0 \end{cases} \Rightarrow \begin{cases} B_0 = \frac{2}{3} R_0 V_0 \\ B_1 = 0 \\ B_2 = -\frac{2}{3} R_0^3 V_0 \end{cases}$$

$$\therefore V(r,\theta) = \frac{2R_0V_0}{3r} - \frac{2R_0^3V_0}{3r^3}P_2(\cos\theta)$$

7.

(a) (i)
$$V = 0$$
 when $v = 0$,

(ii)
$$V = 0$$
 when $y = a$,

(iii)
$$V = V_0$$
 when $x = b$,

(iv)
$$V = V_0$$
 when $x = -b$.

(b)
$$V(x, y) = X(x)Y(y) \implies \frac{1}{X} \frac{\partial^2 X}{\partial x^2} + \frac{1}{Y} \frac{\partial^2 Y}{\partial y^2} = 0$$

$$\begin{cases} \frac{1}{X} \frac{\partial^2 X}{\partial x^2} = k^2 \implies X(x) = Ae^{kx} + Be^{-kx} \\ \frac{1}{Y} \frac{\partial^2 Y}{\partial y^2} = -k^2 \implies Y(y) = C\sin ky + D\cos ky \end{cases}$$

$$V(x, y) = (Ae^{kx} + Be^{-kx})(C\sin ky + D\cos ky)$$

(c)



B.C.(i)
$$\Rightarrow \sin ka = 0$$
, $k = \frac{n\pi}{a}$ $n = 1, 2, 3, ...$
B.C.(ii) $\Rightarrow D = 0$
B.C.(iii) + B.C.(iv) are symmetric, so $A = B \Rightarrow (Ae^{kx} + Be^{-kx}) = A\cosh(kx)$
 $V(x, y) = \sum_{n} C_n \cosh(kx) \sin ky$
B.C.(iv) $V_0 = \sum_{n} C_n \cosh(\frac{n\pi}{a}b) \sin(\frac{n\pi}{a}y)$
 $C_n = \frac{2}{a} \frac{V_0}{\cosh(n\pi b/a)} \int_0^a \sin(\frac{n\pi}{a}y) dy = \frac{2}{a} \frac{V_0}{\cosh(\frac{n\pi}{a}b)} \frac{a}{n\pi} (2) = \frac{4}{n\pi} \frac{V_0}{\cosh(\frac{n\pi}{a}b)}$

$$V(x,y) = \frac{4V_0}{\pi} \sum_{n=1,3,5,...} \frac{1}{n} \frac{\cosh(n\pi x/a)}{\cosh(n\pi b/a)} \sin(n\pi y/a)$$

8.

(a)
$$\mathbf{F}_{e} = q \cdot \mathbf{E} = -q \nabla V$$
 and $V = \frac{1}{4\pi\varepsilon_{0}} \frac{p \cos \theta}{r^{2}}$

$$\nabla V = \frac{\partial V}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{\mathbf{\theta}} + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \hat{\mathbf{\phi}} = \frac{p}{4\pi\varepsilon_{0}} \left[-\frac{2\cos \theta}{r^{3}} \hat{\mathbf{r}} - \frac{\sin \theta}{r^{3}} \hat{\mathbf{\theta}} \right]$$

$$\mathbf{F}_{e} = \frac{qp}{4\pi\varepsilon_{0}} \left[\frac{2\cos \theta}{r^{3}} \hat{\mathbf{r}} + \frac{\sin \theta}{r^{3}} \hat{\mathbf{\theta}} \right]$$

(b)
$$\mathbf{F}_{g} = -mg\hat{\mathbf{z}} = mg(-\cos\theta\hat{\mathbf{r}} + \sin\theta\hat{\mathbf{0}}), \ \mathbf{F}_{e} = \frac{qp}{4\pi\varepsilon_{0}} \left[\frac{2\cos\theta}{r^{3}} \hat{\mathbf{r}} + \frac{\sin\theta}{r^{3}} \hat{\mathbf{0}} \right]$$

$$\mathbf{F}_{total} = \mathbf{F}_{g} + \mathbf{F}_{e} = (-mg + \frac{2qp}{4\pi\varepsilon_{0}r^{3}})\cos\theta\hat{\mathbf{r}} + (mg + \frac{qp}{4\pi\varepsilon_{0}r^{3}})\sin\theta\hat{\mathbf{0}}$$

(c)
$$U_g = mgz = mgr \cos \theta$$
, $U_e = \frac{qp}{4\pi\varepsilon_0} \frac{\cos \theta}{r^2} \implies U_{total} = (mgr + \frac{qp}{4\pi\varepsilon_0 r^2})\cos \theta$

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